

The Molecular and Crystal Structures of Two Diastereomeric 3,4-Dihydropyrrromethenones**

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The molecular and crystal structures of the two racemic diastereomeric 3,4-dihydropyrrromethenone derivatives **1** and **2** of configurations (*Z*) and (*E*) were determined at 93 K (**1**, **2**) and at room temperature (**1**). From an analysis of the crystallographically observed temperature coefficients a pseudorotation flexibility of the pyrrolidinone ring in the crystal is deduced. In both compounds a nearly orthogonal arrangement between the two heterocyclic ring systems is observed, which is ascribed to the steric bulk of the substituents in positions 1 and 3 of the pyrrolidinone ring.

(Keywords: 3,4-Dihydropyrrromethenone; Molecular conformation of 3,4-dihydropyrrromethenones; Phytochrome model compounds; Temperature factor analysis)

Die Struktur zweier diastereomerer 3,4-Dihydropyrrromethenone

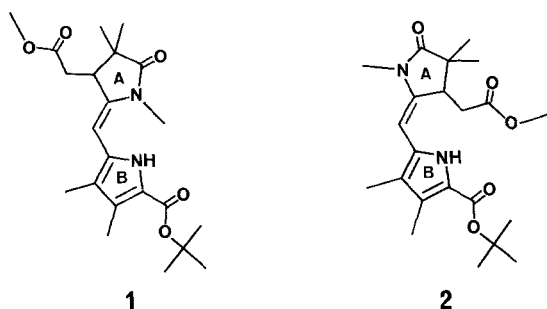
Die Molekül- und Kristallstruktur der beiden racemischen diastereomeren 3,4-Dihydropyrrromethenonderivate **1** und **2** mit Konfigurationen (*Z*) und (*E*) wurden bei 93 K (**1**, **2**) und bei Raumtemperatur (**1**) bestimmt. Aus einer Analyse der kristallographisch beobachteten Temperaturfaktoren wird auf eine Pseudorotations-Flexibilität des Pyrrolidinonringes im Kristall geschlossen. In beiden Verbindungen beobachtet man eine nahezu orthogonale Orientierung der beiden Fünfringe zueinander, die der Raumerfüllung der Substituenten in Positionen 1 und 3 des Pyrrolidinonringes zugeschrieben wird.

Introduction

Investigation of chemical, spectroscopic and structural properties of partial structures is used extensively in the study of linear tetrapyrroles.

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The biochemical significance of this class of compounds with its immense constitutional, conformational and tautomeric variability has attracted increasing attention, particularly since it became evident that the photoreceptor molecule of the plant pigment phytochrome is also a linear tetrapyrrole of the 2,3-dihydrobilatriene type. Phytochrome is responsible for the photomorphogenesis. Several forms of phytochrome are spectroscopically discernible, which are interconverted by light and/or thermal dark reactions. It is highly probable that at least some of these spectroscopic forms are associated with different constitutions, diastereomers or conformations of the phytochromobilin molecule¹.



In the present communication, we report the crystal structures of the two diastereomeric 3,4-dihydropyrromethenones **1** and **2**, both in racemic form. The two molecules differ only with respect to the configuration about the exocyclic double bond, leading to (*Z*) and (*E*) forms (**1** and **2**). Formally, the chromophore of **1** is identical to the “left half” of the chromophore of the phytochromobilin molecule, for which (*Z*)/(*E*) diastereomerism is occasionally discussed as being involved in its biological function. Both compounds were prepared in the course of an extensive series of investigations of differently substituted dihydropyrromethenones², which demonstrated—among other—that the relative thermodynamic stabilities of (*Z*) versus (*E*) forms in dihydropyrromethenones depends on the relative bulk of the substituents in position 3 and on the lactam nitrogen atom. Thus, while the (*Z*)-configuration is thermodynamically favoured for compounds with a proton on the lactam nitrogen, methylation at this position—as in compounds **1** and **2**—inverts the thermodynamic stabilities and makes the (*E*)-diastereomer more stable.

The crystal structures were determined to contribute to the understanding of the above situation. Also, the conformation about the formal exocyclic single bond was of interest: the large substituents on the

pyrrolidine ring made a highly non-coplanar conformation of the two heterocyclic rings probable, and it was of interest to observe the effect of a large twist angle on the length of the (formal) exocyclic single bond.

To our knowledge, only three crystal structures of dihydropyrrromethenones have appeared in the literature so far³; they are the structures of the (*Z*) and (*E*) isomers of 5'-ethoxycarbonyl-3,4-dihydro-3',4'-dimethyl-5(1*H*)-2,2'-pyrrromethenone and the (*E*)-isomer of the same compound with an ethyl group in place of the lactam proton. We will use the results of these structure determinations in the subsequent discussion of our structures. Another structural investigation of a (*Z*)-derivate was announced in a preliminary communication⁴.

Both structure determinations reported in this communication were carried out at low temperature (93 K) to attain maximum accuracy. The structure of **1** was also determined at room temperature; the combination of room temperature and low temperature data permits a test on the presence or absence of static disorder in the crystal. Finally, we will present the results of an analysis of the temperature factors for the two low temperature structures to obtain an indication about the flexibility of the molecules in the crystal.

Results and Discussion

Tables 1 and 2 list atomic coordinates and equivalent isotropic temperature coefficients with associated standard deviations for the room temperature structure of **1** and for the two low-temperature structures, for which anisotropic temperature coefficients are given in Table 3. Bond lengths and bond angles calculated from the low-temperature coordinates are summarized in Fig. 1, which also defines the atom numbering used.

Molecular Constitution and Conformation

The structure analyses (Figs. 2 and 3) confirm the (*Z*)- and (*E*)-configurations for **1** and **2**, respectively. As expected, both structures display pronounced non-planarity between the two five membered rings; the large values observed for the twist angle about the $C_{\alpha}-C_m$ single bond (C4-C5-C6-N11, Fig. 1) make a classification of the conformation in terms of "*syn*" or "*anti*" somewhat pointless; however, since both torsion angles have absolute values beyond 90°, the two molecules should most appropriately be termed (*Z*)- and (*E*)-anticlinal, respectively.

The main factor responsible for the anticlinal conformation is probably steric repulsion between the substituents on the pyrrole and the pyrrolidine rings, which prohibits a periplanar conformation. The exact value of the C5-C6 dihedral angle, however, may well be significantly

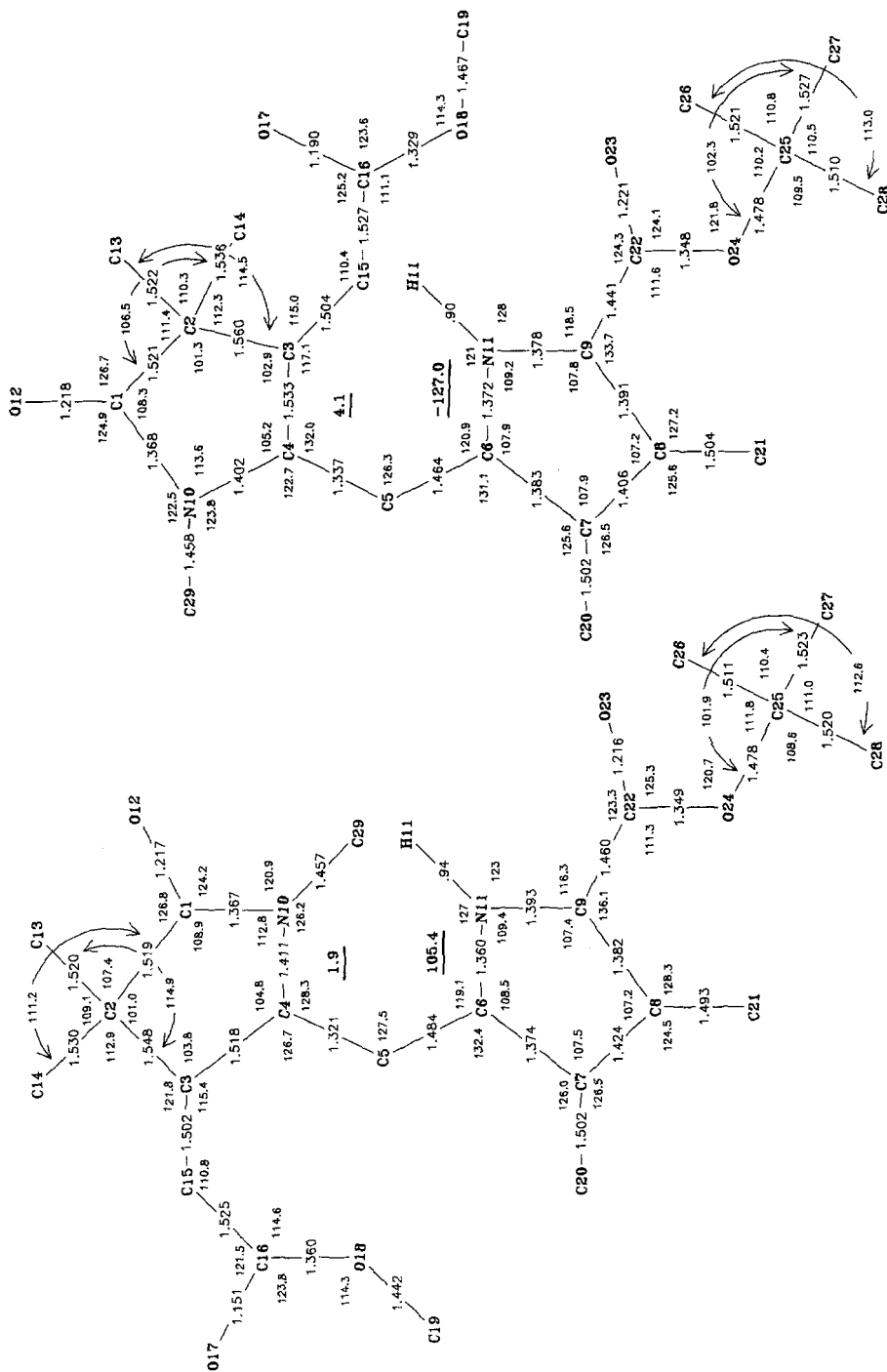


Fig. 1. Calculated bond lengths and bond angles for the low temperature crystal structures of **1** (left) and **2** (right). Two torsion angles about the two $C^m - C^n$ bonds are also given (underlined). Derived standard deviations are as follows: **1**: $X - Y$, $\sigma \sim 0.005 - 0.001$; $N - H$, $\sigma \sim 0.05$ Å; $X - Y - Z$, $\sigma \sim 0.4 - 0.6$; $X - N - H$, $\sigma \sim 2.8^\circ$; **2**: $X - Y$, $\sigma \sim 0.003 - 0.005$; $N - H$, $\sigma \sim 0.2 - 0.3^\circ$; $X - N - H$, $\sigma \sim 2.0^\circ$ ($X, Y, Z = C, N, O$).

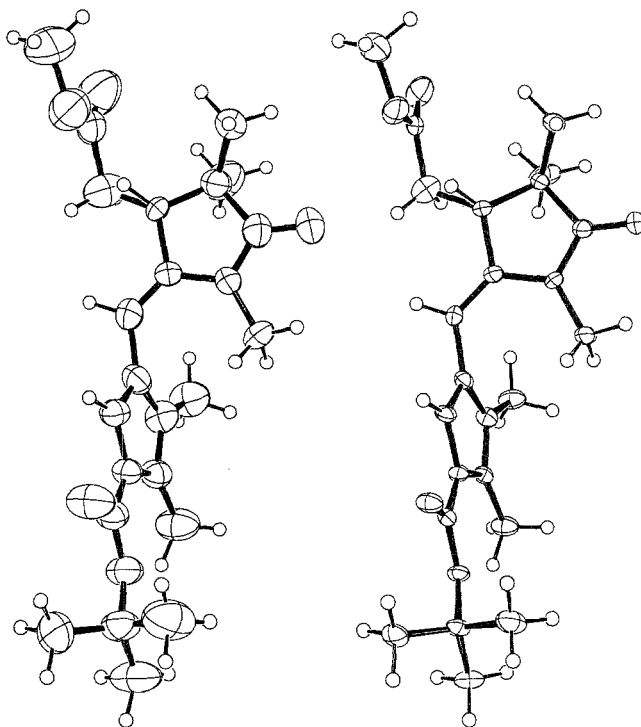


Fig. 2. ORTEP drawing of the room temperature (left) and low temperature crystal structures (right) of **1**. Temperature ellipsoids of C, N and O are drawn at the 50% level; H atoms are represented as spheres of radius 0.1 Å

influenced by intermolecular hydrogen bonding (see below) and dispersion forces, whose collective effect is difficult to assess.

The degree of conjugation across the formal $C_m - C_\alpha$ "single" bond is frequently discussed in crystal structures of linear di- and tetrapyrroles. Fig. 4 shows a plot of the C 5–C 6 bondlength versus the absolute value of the dihedral angle about this bond for the three structure determinations on **1** and **2** and for the three dihydropyromethenone crystal structures described in Ref.³. In spite of the limited data, there clearly is a positive correlation between dihedral angle and bond length, which varies between 1.44 and 1.50 Å for $\tau = 0$ and 90° , respectively. It is well established that the conjugation across the "single" bond is more pronounced in pyromethenones⁵, whose corresponding bond length is observed at about 1.43 Å for the coplanar conformation.

As observed in previously reported dihydropyromethenone structures³, the ester carbonyl group is approximately coplanar with the

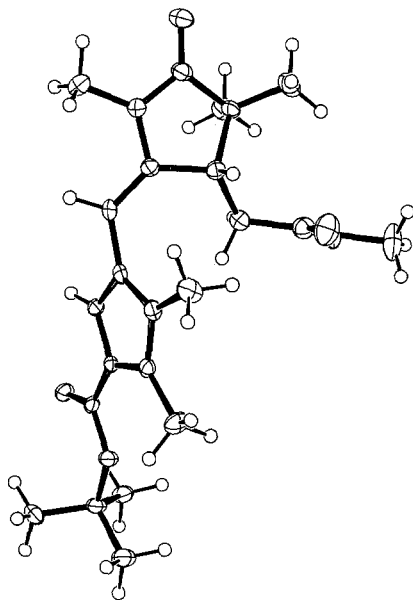


Fig. 3. ORTEP drawing of **2** as observed in its low temperature crystal structure

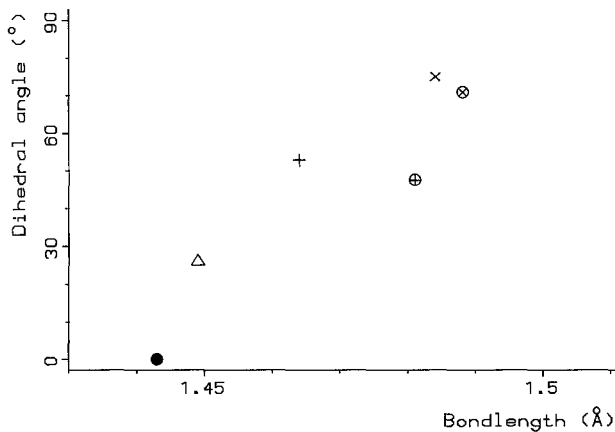


Fig. 4. Plot of the $C_5 - C_6$ bond length versus the absolute value of the smaller of the two torsion angles for 6 crystal structure determinations of dihydropyrromethenones. × **1**, low temperature; ⊗ **1**, room temperature, + **2**, low temperature; ⊕ compound I, ● compound II, Δ compound IV from Ref.³

pyrrole ring (torsion angles C 8–C 9–C 22–O 24: -8.5° in **1** and 4.4° in **2**), with the carbonyl oxygen syn to the pyrrole nitrogen.

The lower free enthalpy of **2** as opposed to **1** has been ascribed to the effect of steric repulsion between the methyl group on the lactam nitrogen and the pyrrole ring². Indeed, there is crystallographic evidence in support of this explanation: both, the C 29–N 10–C 4 and N 10–C 4–C 5 angles are significantly larger in **1** than in **2** (by 2.4° and 5.6° , respectively); these are the largest differences between corresponding angles in the two structures.

Pyrrolidine Ring Conformation

While the pyrrole ring is planar within experimental error in both compounds, the pyrrolidine rings show significant deviations from planarity. The conformation of five-membered rings is conveniently described in terms of a puckering amplitude q and a phase angle θ , using the equation for a pseudorotating pentagon⁶

$$Z_j = (2/5)^{1/2} \cdot q \cdot \cos(4\pi j/5 + \theta)$$

(Z_j is the normal deviation of atom j from a least-squares plane through the five-membered ring). The values observed in the two structures are $q = 0.291$ (**1**) and 0.294 \AA (**2**) and $\theta(\text{mod } 36) = 28.2$ (**1**) and 23.3° (**2**), respectively.

Since $\theta(\text{mod } 36)$ —values of 0° or 36° correspond to envelope and of 18° to twist conformations, the pyrrolidine rings are intermediate between envelope and twist forms in both structures.

The puckering amplitudes are remarkably large in both compounds, compared to values between 0.06 and 0.11 \AA observed in other dihydropyromethenone crystal structures⁵. We searched the Cambridge Crystallographic Data File⁷ for pyrrolidine rings with exocyclic double bonds in positions 2 and 5 and obtained 18 structures containing such a fragment. The average value of q for those 18 rings was 0.093 \AA , the largest value was 0.23 \AA . We ascribe the large value of q in **1** and **2** to the effect of the large substituents on C 2 and C 3, which appear to enforce a more or less staggered conformation between the methyl groups on C 2 and the acetic acid ester group on C 3 (torsion angles C 13–C 2–C 3–C 15: -45.1° in **1** and -83.2° in **2**; C 14–C 2–C 3–C 15: 80.9° in **1** and 43.6° in **2**).

Packing

Both compounds have an “acidic” hydrogen (H 11 on the pyrrole nitrogen atom N 1) capable of forming an intermolecular H-bond to any of the carbonyl oxygen atoms. As expected, this H-bond has a decisive influence on the packing of both structures (Fig. 6).

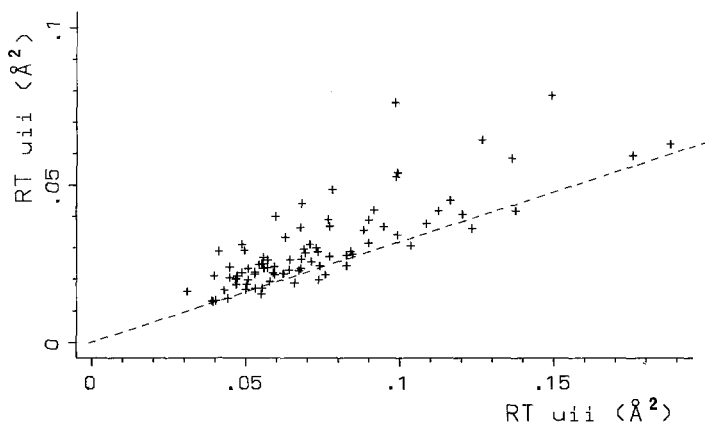


Fig. 5. Scatterplot of room temperature u_{ii} values (abscissa) versus the corresponding low temperature values for the crystal structure of **1**

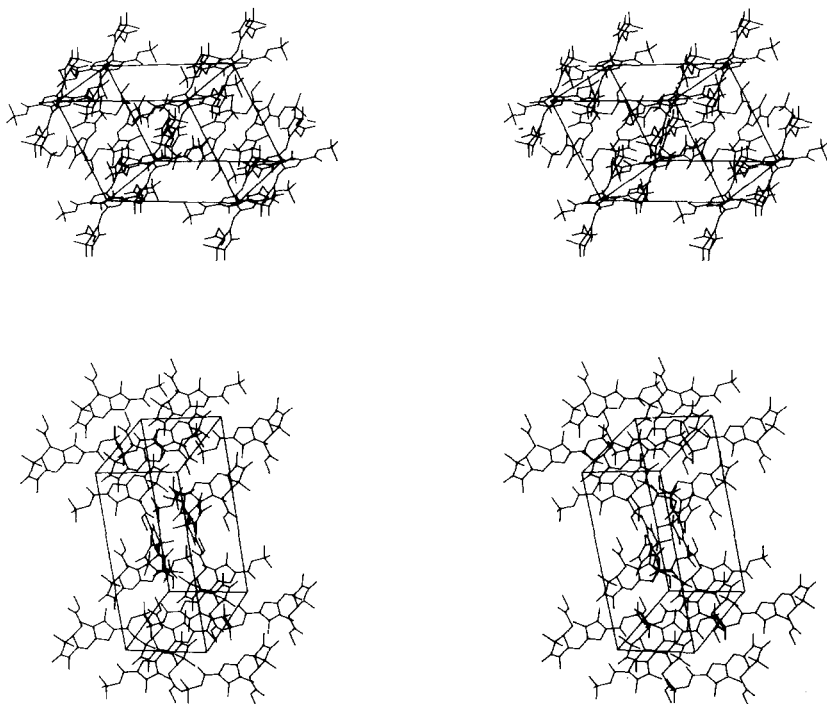


Fig. 6. Stereoscopic packing diagrams for the crystal structures of **1** (top) and **2** (bottom). The crystallographic axes are oriented as follows: **1**: X towards the viewer, Y left to right, Z upwards; **2**: X upwards, Y towards the viewer, Z left to right

In **1** the H-bond involves the lactam oxygen atom of ring A (H 11 – O 12_{-1/2-x, -1/2+y, -z}, $d = 1.95 \text{ \AA}$) of a molecule related by the operation of the 2_1 screw axis, leading to the formation of one-dimensional chains of molecules running along the monoclinic axis. The packing of **2**, on the other hand, is dominated by the formation of centrosymmetric dimers involving the carbonyl oxygen of the *t*-butylester grouping (H 11 – O 23_{-x, -y, 2-z}, $d = 1.99 \text{ \AA}$).

No pronounced base-stacking is observed in either one of the two structures.

Comparison of Room Temperature and Low Temperature Structures of 1

Fig. 2 shows ORTEP-drawings of the asymmetric unit of **1** as observed at the two temperatures; a scatterplot of low-temperature (LT) versus room-temperature (RT) u_{ii} values for the anisotropically refined atoms (i.e. C, N, O) is shown in Fig. 5. To a first approximation (each atom vibrating independently in a harmonic potential^{6,8,9}), such a plot should yield a straight line through the origin with slope T_1/T_2 (T is the absolute temperature). This line (with slope $93/293 \approx 0.32$) is also indicated in the plot, and it is evident that the LT u_{ii} values are systematically higher than predicted by the model. Barring experimental errors, the discrepancy is most probably due to inherent limitations of the harmonic approximation, in particular the effect of zero-point vibrations. Nevertheless, the plot gives little indication for disorder in the structure.

As expected, RT bond lengths are systematically shorter [by 0.005 (1) \AA] than LT ones. The difference is readily explained as the result of anisotropic libration.

Analysis of Vibration Ellipsoids

Energies involved in bond stretching are large compared to those associated with other deformations in a crystal (lattice vibrations, torsion and bond angle deformations). Therefore, mean square vibrational amplitudes of atoms connected by a chemical bond should be approximately equal along the bond direction¹¹. This property, which should also hold for any two atoms belonging to a rigid group, can be used to test the applicability of the rigid-body model to a molecule or molecular group¹²: with Z_{kl}^2 being the mean-square vibration amplitude of atom k along the line connecting atoms k and l , the test calls for computation of the $n \times n$ matrix of $\Delta_{kl} = Z_{kl}^2 - Z_{lk}^2$ values ($n =$ number of atoms), and accepts the rigid-body assumption if all the Δ -values are zero within experimental error.

Δ -values differing significantly from zero indicate non-rigidity. U_{ij} values from structure refinement against X-ray data are prone to

systematic errors and tend to absorb charge-density deformation contributions. However, the fact that Δ -values along chemical bonds should always be zero (see above) can serve as an indication for the usability of the u_{ij} values at hand.

Instead of Δ_{kl} , we will in the following use $\delta_{kl} = \Delta_{kl}/\sigma(\Delta_{kl})$, thereby replacing differences in mean square vibrational amplitudes by their statistical significance.

In the course of an analysis of the temperature factors for molecular flexibility, two types of question may be encountered: (1) does a fragment behave as a rigid group and (2) does one rigid fragment “vibrate” relative to a second rigid fragment? Within the inherent limits of this approach, both types of questions can be treated by inspection of the appropriate elements of the δ matrix. To check whether all elements of the corresponding subset of the δ_{kl} matrix are zero within experimental error, we use the well-known technique of the half normal probability plot (HNPP)⁹, which amounts to a plot of the δ_{kl} versus their statistical expectation values. If the hypothesis of zero δ_{kl} is fulfilled, the plot yields a straight line through the origin with unit slope. A HNPP with slope < 1 indicates δ_{kl} systematically different from zero or underestimation of the $\sigma(\Delta_{kl})$.

In the present application, Δ_{kl} were computed from the U_{ij} derived from the least squares refinement. Their associated standard deviations $\sigma(\Delta_{kl})$ were obtained from the $\sigma(U_{ij})$, neglecting covariant terms.

Flexibility of 1 and 2 in the Crystal

A first rough indication of the most flexible regions can be obtained from the equivalent isotropic temperature factors listed in Table 1: in all three structures, the largest u_{iso} values are found for the atoms of the methyl ester and for the ring A substituents, while the atoms of ring B and the *t*-butylester show u_{iso} values below average.

The HNPP of δ_{kl} along chemical bonds is shown in Fig. 7 A for the low-temperature data of **1**. Since δ_{kl} along bond directions should be zero, the figure is no more than a test for the quality of the data. Evidently, the plot shows a slope less than unity (~ 0.62), which will have to be born in mind for the interpretation of the following HNPP's.

Fig. 7 B shows a HNPP of the δ_{kl} for all intra-fragment vectors involving the pyrrole ring atoms and the atoms directly attached to it. Since δ_{kl} values are insensitive to out-of-plane vibrations of planar groups⁹, the observed slope of ~ 0.95 is as expected and yields another test for the quality of the u_{ij} values.

The analogous HNPP for the pyrrolidine ring (Fig. 7 C) is strikingly different: with an observed slope less than 0.3, we can safely conclude that

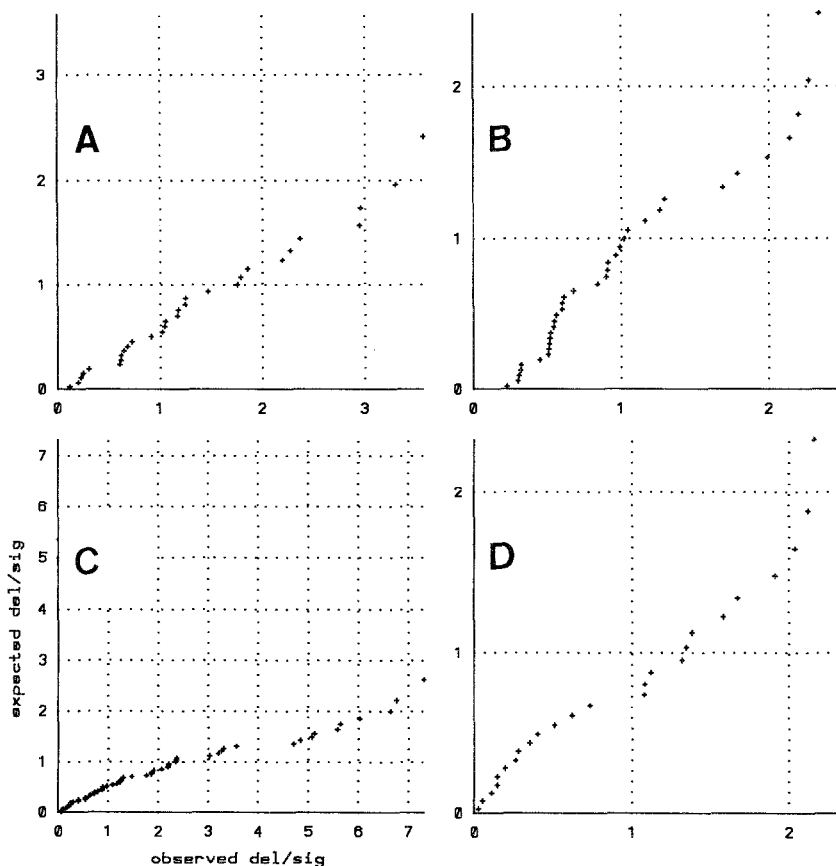


Fig. 7. Half-normal probability plots of δ_{kl} from the low-temperature structure of **1** (see text): A: k, l = all pairs of bonded atoms; B: k, l = atoms 1–5, 10, 12–15, 29; C: k, l = 5–9, 11, 20–22; D: k = 6, 9, 11, 28, 29; l = 1, 2, 3, 4, 10

the ring shows “flexibility” in the crystal, i.e. the observed u_{ij} are not compatible with a rigidly librating molecular fragment. Non-planar five-membered rings are known to show pseudorotation, which makes the above result quite plausible.

A more unexpected result emerges from Fig. 9 D, which shows a HNPP of the δ_{kl} for vectors between pyrrole and pyrrolidine ring. Since these two fragments are connected by a formal single bond, one might anticipate that they librate more or less independently. However, the HNPP yields no experimental evidence to reject a rigid-molecule

Table 1. Atomic coordinates and equivalent isotropic temperature factors ($\cdot 10^4$, U -values in \AA^2) for the non-hydrogen atoms in the crystal structures of **1** (A: room temperature, B: low temperature), **2** (C: low temperature). Estimated standard deviations are given in italics. The equivalent isotropic temperature factors were obtained as one third of the trace of the orthogonalized U_{ij} tensor, and it corresponds to an isotropic temperature factor of the form $T = \exp(-8\pi^2 \cdot u \cdot \sin^2 \theta / \lambda^2)$

Atom	A			B			C																	
	X/a	Y/b	Z/c	U _{iso}	X/a	Y/b	Z/c	U _{iso}	X/a	Y/b	Z/c	U _{iso}												
C1	-3060	7	3549	6	-1006	5	628	60	-3088	4	3545	4	-995	3	243	31	4271	3	5186	3	8808	2	245	17
C2	-3188	6	2884	6	-1814	4	538	54	-3223	4	2876	4	-1824	3	267	32	-4117	2	4497	3	8130	2	224	16
C3	-2572	7	1815	6	-1395	4	776	60	-2564	5	1790	6	-1407	4	458	41	-2781	2	4387	3	8232	1	209	16
C4	-1713	7	2172	6	-517	4	478	53	-1687	4	2157	4	-521	3	251	31	-2424	2	4297	3	9027	1	181	15
C5	-770	7	1649	6	-25	5	600	60	-702	5	1634	5	-47	3	256	32	-1485	3	3795	3	9433	1	199	16
C6	89	6	1916	6	876	5	537	57	115	4	1911	4	862	3	212	29	-468	2	3241	3	9198	1	184	15
C7	905	7	2731	6	1190	5	580	57	950	4	2730	4	1182	3	228	30	255	2	3813	3	8791	1	202	15
C8	1429	6	2583	6	2120	4	514	55	1488	4	2576	4	2118	3	185	28	1122	2	2853	3	8739	1	189	15
C9	898	6	1685	6	2394	4	452	53	943	4	1672	4	2340	3	203	28	918	2	1707	3	9126	1	160	15
N10	-2190	5	3131	5	-308	3	558	45	-2187	3	3117	3	-294	3	229	25	-3308	2	4946	3	9296	1	218	14
N11	75	5	1286	5	1570	4	522	45	106	3	1271	4	1556	3	219	26	-55	2	1961	3	9402	1	178	13
O12	-3624	4	4370	4	-974	3	782	37	-3672	3	4358	3	-950	2	280	21	-5116	2	5808	2	8915	1	365	14
C13	-2532	11	3547	7	-2309	6	983	82	-2622	5	3529	7	-2393	4	454	44	-4663	3	5337	4	7499	2	347	20
C14	-4511	7	2740	7	-2392	6	853	62	-4531	5	2701	5	-2393	4	355	35	-4697	3	3064	4	6120	2	320	20
C15	-2119	10	1030	9	-1688	6	1122	82	-2123	7	982	7	-1927	5	592	51	-2389	3	3257	3	7807	2	244	17
C16	-3098	9	566	7	-2668	6	809	77	-3136	5	546	5	-2728	4	399	38	-2435	3	3749	4	7050	2	283	18
O17	-3208	7	841	5	-3380	4	1282	64	-3264	4	829	3	-3440	3	505	31	-2119	2	4827	3	6868	1	406	15
O18	-3792	7	-200	5	-2521	4	1048	55	-3810	4	-247	3	-2538	2	476	29	-2884	2	2765	2	6612	1	361	14
C19	-4603	11	-732	11	-3317	8	1325	91	-4634	6	-793	6	-3301	5	466	44	-2992	4	3129	5	5873	2	484	28
C20	1163	7	3611	7	647	5	791	73	1217	5	3641	5	844	3	321	35	118	3	5209	4	8458	2	310	20
C21	2363	7	3287	7	2700	5	826	70	2441	5	3300	5	2719	4	315	34	2085	3	3050	4	8338	2	293	19
C22	930	7	1166	7	3148	5	534	59	972	4	1136	4	3161	3	217	30	1477	2	395	3	9259	1	182	15
O23	277	5	482	5	3187	3	810	45	268	3	427	3	3188	2	277	23	1125	2	-534	2	9586	1	232	12
O24	1775	4	1612	4	3947	3	618	38	1638	3	1563	3	3874	2	216	20	2435	2	317	2	8977	1	204	11
C25	1865	8	1323	7	4766	4	664	67	1829	4	1255	4	4786	3	228	31	3130	2	-951	3	9015	2	218	16
C26	2098	10	119	8	4913	6	911	84	2067	5	20	5	4935	4	324	38	3620	3	-1294	4	9771	2	268	18
C27	2987	9	1983	11	5304	5	985	94	3060	5	1853	6	5347	4	324	37	4085	3	-544	4	8621	2	291	19
C28	790	10	1735	14	4911	6	1137	113	852	5	1715	7	4858	4	394	42	2413	3	-2097	3	8639	2	248	18
C29	-1902	7	3629	6	556	4	641	56	-1937	4	3623	5	584	3	271	32	-3224	3	5335	4	10027	2	307	20

Table 2. Atomic coordinates and isotropic temperature factors ($\cdot 10^3$) for the hydrogen atoms in the low temperature crystal structures of **1** (A) and **2** (B). Note that standard deviations in the coordinates may be underestimated due to distance constraints during the refinement (see text in experimental section)

A					B											
Atom	X/a	Y/b	Z/c	U _{iso}	X/a	Y/b	Z/c	U _{iso}								
H -C3	-294	5	101	3	-159	4	0	24	-245	2	526	1	808	1	18	8
H -C5	-52	3	99	4	-29	3	0	11	-147	2	381	3	993	2	22	8
H -N11	-32	4	61	4	153	3	22	14	-36	3	145	3	971	2	28	9
H1-C13	-264	5	308	4	-288	2	54	17	-449	3	493	4	706	1	55	12
H2-C13	-176	2	371	5	-201	3	73	20	-432	2	628	1	766	2	33	9
H3-C13	-303	5	428	2	-249	4	58	21	-550	1	543	3	751	2	43	10
H1-C14	-500	5	234	5	-205	3	65	20	-465	3	261	3	767	1	51	11
H2-C14	-461	4	224	4	-294	2	40	16	-552	1	324	3	816	2	35	9
H3-C14	-489	4	346	2	-258	3	35	16	-430	2	248	3	851	1	34	9
H1-C15	-159	8	160	6	-200	7	211	53	-150	1	311	3	798	2	41	10
H2-C15	-182	9	38	6	-148	6	229	50	-272	2	237	1	783	1	10	7
H1-C19	-420	5	-125	4	-362	4	68	22	-222	1	340	3	577	2	33	10
H2-C19	-508	4	-128	4	-302	3	50	18	-335	3	406	2	581	2	76	15
H3-C19	-526	4	-25	4	-366	4	94	25	-335	3	230	3	562	2	74	14
H1-C20	69	3	357	4	0	1	30	14	89	1	564	3	848	2	31	9
H2-C20	102	5	437	2	87	3	47	17	-29	3	515	4	796	1	68	13
H3-C20	209	1	375	4	75	3	32	14	-30	3	587	3	872	2	57	12
H1-C21	214	4	407	2	279	3	44	16	196	3	253	3	789	1	33	9
H2-C21	283	5	301	5	334	2	69	20	284	2	271	4	860	2	68	13
H3-C21	313	3	347	4	252	3	55	17	223	3	403	1	824	2	57	12
H1-C26	233	5	-13	5	560	1	53	19	422	2	-203	2	981	1	25	8
H2-C26	269	3	-25	3	469	3	12	12	303	2	-157	3	1005	1	30	9
H3-C26	130	3	-39	4	462	3	47	19	402	2	-46	2	999	1	25	8
H1-C27	317	4	168	4	599	1	35	15	462	2	-134	2	860	1	13	7
H2-C27	302	4	287	1	520	3	26	15	456	2	23	2	866	1	36	10
H3-C27	373	3	159	4	515	3	27	15	374	2	-23	3	814	1	34	9
H1-C28	90	5	153	5	558	1	53	18	290	2	-293	2	861	1	26	9
H2-C28	8	2	139	4	455	3	43	16	206	3	-174	3	817	1	47	11
H3-C28	78	5	255	1	489	4	39	19	181	2	-235	3	892	1	16	7
H1-C29	-172	4	307	3	108	2	35	16	-251	2	590	3	1013	2	44	10
H2-C29	-114	3	414	3	72	3	49	17	-338	3	450	2	1030	2	71	14
H3-C29	-252	3	410	3	59	3	27	13	-388	2	597	3	1006	2	54	11

hypothesis. Although this cannot be taken as a proof for rigidity, it does make the existence of large relative vibrations between the two molecular fragments improbable.

We refrain from giving the analogous analysis for the u_{ij} 's of structure **2**, since essentially the same result is obtained: again, there is evidence for internal flexibility of the pyrrolidine ring, while no indication for vibrations between pyrrole ring and the planar part of the pyrrolidine ring is obtained.

Table 3. Anisotropic temperature coefficients ($\cdot 10^4$, in \AA^2) for non-H atoms in the low-temperature crystal structures of **1** (A) and **2** (B). The temperature factor has the form: $T = \exp(-2\pi^2(h^2 a^{*2} u_{11} + \dots + 2hka^*b^*u_{12} + \dots))$

Atom	A						B					
	u11	u22	u33	u12	u13	u23	u11	u22	u33	u12	u13	u23
C1	197 25	291 32	217 28	1 25	48 22	-41 25	207 15	256 18	283 17	48 13	73 13	80 15
C2	210 27	334 33	201 28	98 25	2 22	-52 26	129 13	320 18	219 16	12 14	21 12	44 14
C3	421 36	442 45	288 31	235 35	-153 37	-206 32	178 14	245 17	206 16	-21 13	42 13	15 14
C4	247 27	311 32	161 27	53 26	28 22	-62 24	157 13	194 15	198 15	-17 12	44 12	13 13
C5	295 29	233 31	215 28	79 26	64 23	-28 26	238 15	225 17	129 14	4 13	21 12	34 13
C6	171 25	248 29	200 26	39 23	46 21	13 23	186 14	214 16	134 14	4 13	-23 11	4 13
C7	191 26	245 30	238 28	-29 24	62 22	29 25	191 14	228 16	178 15	-24 13	10 12	26 13
C8	132 22	219 30	197 26	0 22	50 20	11 23	158 13	255 17	148 14	-42 13	10 11	-7 13
C9	128 23	235 29	209 27	-18 23	11 21	-31 24	138 13	223 16	118 14	-14 12	18 11	-5 12
N10	236 23	260 25	185 22	36 21	37 19	-74 20	224 13	236 14	195 13	36 11	44 11	13 11
N11	152 22	238 26	205 23	-16 19	-14 18	9 21	154 11	224 14	162 13	-11 11	45 10	42 11
O12	276 19	262 21	286 19	67 17	80 16	-24 17	274 11	454 15	377 14	146 11	89 11	80 12
C13	362 35	762 53	283 32	-71 38	168 29	-166 36	209 15	337 24	337 20	21 17	-10 15	116 19
C14	229 28	311 35	369 35	37 27	-60 26	-77 30	221 16	337 21	403 21	-44 16	62 16	30 18
C15	644 49	528 50	485 42	46 43	47 38	-154 40	235 16	252 18	244 16	25 14	39 13	16 14
C16	418 36	187 32	400 36	95 28	-119 30	-71 29	278 17	283 19	284 18	24 15	40 14	-27 16
O17	594 29	377 27	539 29	-66 23	193 25	-89 24	539 16	383 14	359 14	-71 13	117 12	41 12
O18	786 31	368 26	213 20	99 24	99 21	8 19	505 15	283 13	224 12	31 12	-46 11	8 11
C19	341 36	408 42	584 44	-93 33	82 34	88 37	821 33	404 27	199 19	-34 25	19 21	20 18
C20	367 32	390 36	188 27	-88 31	75 25	34 28	285 17	305 20	308 21	-9 16	61 16	93 17
C21	271 29	356 36	263 29	-104 28	27 24	17 28	271 17	306 20	340 20	-49 16	156 15	50 17
C22	220 26	222 30	183 26	61 25	40 22	18 24	153 13	272 17	123 13	-29 13	26 11	-22 13
O23	288 21	279 22	217 19	-74 18	29 17	51 17	218 11	266 12	233 11	24 10	98 9	87 10
O24	255 19	213 20	140 17	-23 17	18 15	35 16	177 10	214 11	236 11	44 9	77 9	46 9
C25	242 28	299 33	131 25	17 24	51 22	56 23	161 14	235 17	204 16	43 13	58 13	26 14
C26	419 37	268 35	241 32	-34 30	63 29	93 27	204 16	336 20	246 17	28 15	-11 14	44 16
C27	307 32	452 42	170 29	-85 32	30 26	-4 30	207 16	328 21	356 20	34 16	94 15	41 17
C28	315 34	630 51	227 32	118 36	82 27	25 34	214 15	279 19	269 18	16 15	79 14	-49 16
C29	233 27	365 35	193 26	54 26	34 22	-43 25	325 18	359 22	260 18	52 17	113 15	-14 16

Table 4. *Summary of experimental conditions*

	1 (low temp.)	1 (room temp.)	2 (low temp.)	
chemical formula	$C_{22}H_{32}N_2O_5$		$C_{22}H_{32}N_2O_5$	
crystal size [mm]	.1 × .25 × .3	.15 × .3 × .3	.2 × .25 × .45	
temperature [K] (cold stream)	93 ± 1	293 ± 2	93 ± 1	
space group	P 2 ₁ /a			
cell dimensions	a [Å]	12.019 (4)	12.132 (3)	11.830 (2)
	b [Å]	12.073 (4)	12.228 (4)	9.739 (1)
	c [Å]	16.070 (3)	16.250 (4)	19.483 (4)
	β [°]	110.61 (3)	111.08 (2)	100.28 (1)
	V [Å ³]	2182.6	2249.2	2208.8
number of molecules per unit cell	4			
calculated density [g cm ⁻³]	1.231	1.194	1.216	
number and θ-range [°] of reflections used to refine cell constants	15 24 ≤ 2θ ≤ 30	14 14 ≤ 2θ ≤ 32	23 19 ≤ 2θ ≤ 21	
limits for data collection [°]	0 ≤ 2θ ≤ 60 - ≤ 16 h ≤ 16 0 ≤ k ≤ 16 0 ≤ l ≤ 22	0 ≤ 2θ ≤ 50 0 ≤ h ≤ 14 0 ≤ k ≤ 14 - 19 ≤ l ≤ 18	0 ≤ 2θ ≤ 60 - 18 ≤ h ≤ 18 0 ≤ k ≤ 15 0 ≤ l ≤ 27	
scan width [°] (ω/θ-scan)	Δω = 1.2	Δω = 1.2	Δω = 1.2	
number/frequency of standard refs.	3/100	3/100	3/100	
max. variation of standard intensities	± 2%	± 3%	± 3%	
number of observed reflections	7061	4525	7214	
number of independent reflections	6357	3973	6451	
number of reflections with F _o > 4σ(F _o)	1994	1218	2675	
R/R _w -factors	0.070/0.067	0.063/0.062	0.055/0.052	
number of parameters	391	294	391	
number of observations (including distance constraints)	2024	1218	2705	
coefficients a/b of weighting system w _i = a/(σ(F _i) ² + bF _i ²)	2.40/0.00035	2.13/0.0004	1.69/0.00007	
highest peak/lowest trough in final ΔF-Fourier synthesis [e Å ⁻³]	0.63/- 0.46	0.27/- 0.21	0.30/- 0.29	

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Experimental

Colourless crystals of **1** and **2** were obtained by dissolving a small amount of material² in about 1 ml chloroform. The solution was placed inside a dessiccator

containing a second open beaker with *n*-hexane and kept in the coldroom for several weeks. The crystals used for structure analysis had to be cut with a razor blade from larger lumps containing several crystals.

All X-ray work was carried out on an extensively modified STOE 4-circle diffractometer (MoK α -radiation, graphite monochromator, $\lambda = 0.71069 \text{ \AA}$) equipped with a NONIUS low-temperature attachment. To prevent deposition of ice on the crystal, the diffractometer was placed inside a glove box. Conditions of data collection and structure refinement are summarized in Table 4.

Data processing involved the usual LP-corrections, but neither absorption- nor extinction correction [$(\mu(\text{MoK}\alpha) = 0.94 \text{ cm}^{-1}$ for **1** and 0.93 cm^{-1} for **2**)]. The two structures were solved by direct methods and refined first by full matrix and subsequently by blocked least-squares techniques. In view of their superior quality, the refinement was first carried out for the two low-temperature data sets, applying the following conditions: all non-hydrogen atoms were refined with anisotropic temperature factors, H-atoms were included at observed positions and refined with isotropic temperature factors. In the terminal refinement cycles, aliphatic C–H distances were constrained to assume the same value; this “average” aliphatic C–H distance converged to 1.00 \AA in both low-temperature refinements.

Subsequently, the low-temperature coordinates were used as a starting point for the refinement of **1** against the room-temperature data. Again, all non-H atoms were refined with anisotropic temperature factors, whereas hydrogen atoms were kept at the positions of the low-temperature refinement. An isotropic temperature factor was refined for each H-atom.

Atomic scattering factors were from the International Tables¹³. Computer programs used (in locally modified versions) are summarized in Ref.¹⁴.

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